

Plan: ① Intro/Ising magnets ② Heisenberg systems ③ Quasi Spin Liquids

~~Mott~~

Frustrated Magnets

- Magnets \rightarrow I will focus on nearly localized e^-
 \approx atomic eigenstates w/ weak overlap

usually Mott Insulators

- Get magnetism when partially filled shells

e.g.

Mn^{2+}



$s = 5/2$

lots of complexity + variation due to many different atomic conditions

- How do you know?

- Usually good insulator $\chi_p \sim 1/T$
- Curie Law at high-Temp

$$\chi \sim \frac{A}{T}$$

$$A = \frac{Ng^2 \mu_B^2 S(S+1)}{3k_B} \quad \text{Curie Constant}$$

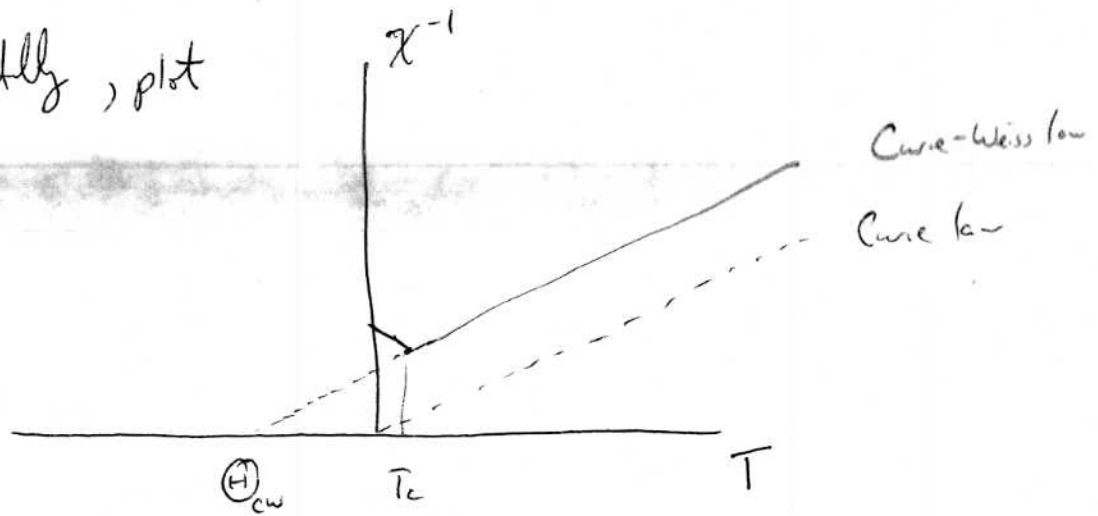
signature of free moments.

- Interactions between spins e.g. $H = \frac{1}{2} \sum_{ij} J_{ij} \vec{S}_i \cdot \vec{S}_j$

(note I choose $J > 0$ for AF)

- AF interactions \rightarrow generally suppresses χ .

Empirically, plot



$$\chi \sim \frac{A}{T - \Theta_{cw}} \quad \Theta_{cw} < 0 \quad \text{for AFs}$$

- This form is obtained in MFT close T_c
- Also in high- T expansion

$$\chi^{-1} \sim A^{-1} (T - \Theta_{cw} - O(1/T))$$

$$\Theta_{cw} = - \frac{(\sum_j J_{ij}) S(S+1)}{3k_B}$$

- $|\Theta_{cw}|$ Gives a measure of strength of negative interactions

- For ~~unfrustrated~~ e.g. cubic lattice AF,

$$|\Theta_{cw}| = T_c^{AF}$$

• "Frustration Fingerprint"

$$f = \frac{|\Theta_{cw}|}{T_c} \gg 1 \longrightarrow \text{ordering is anomalously suppressed. "frustrated"}$$

empirically $f \gtrsim 5-10$ usually accepted as frustrated.

Why $f \gg 1$?

- Frustration \approx competing interactions

→ Cannot satisfy all $J_{ij} \vec{S}_i \cdot \vec{S}_j$ terms simultaneously

→ Get many almost equally good compromises

= *degeneracy*

→ System fluctuates between these degenerate states instead of ordering ← Quark/Neut?

Before getting to any details, talk about why this is interesting

• Emergent low energy scale $T_c \ll \Theta_{CW}$

- Expect some new (universal?) behavior $T_c < T < \Theta_{CW}$?
- Here spins are correlated but not ordered. "spin liquid"

• Quasi-degeneracy

- Perturbations that break degeneracy are only ^{intrinsic} energy scale

→ Similar to FQHE → project T_c into LLL

or

- Very sensitive

- relatively small effects can control G.S.
- controllable?

- Intensity states arise
 - non-collinear magnetic order (c.f. Mastroy))
 - non-magnetic states: dimers, spin-liquids
 - reduced dimensionality

Ising Magnet
theory $\vec{S}_i \rightarrow \sigma_i^z \equiv \sigma_i \quad \sigma_i = \pm 1$ (

practice: need strong ~~single-ion anisotropy~~ single-ion anisotropy
Crystal field + spin-orbit effects

most ~~common~~ ^{Ising-like} \rightarrow f e^- 's in rare earth elements.

but also some d e^- 's in asymmetric environments

~~Difficulties are of Ising-like~~

Difficulties

- Ising axes often not collinear
 $\vec{S}_i \rightarrow \hat{n}_i \sigma_i$ \hat{n}_i : different for different sites
- For f e^- 's, dipolar interactions important.

(both effects are present in famous "spin-ice")
materials I will discuss

NN S_i^z AFs

$$H = J \sum_{\langle ij \rangle} \sigma_i \cdot \sigma_j$$

Note: this is purely classical problem since $[\sigma_i^z, \sigma_j^z] = 0$.

QM's arises if one includes transverse operators, $e.g.$

$$J_{\perp} (S_i^x S_j^x + S_i^y S_j^y)$$

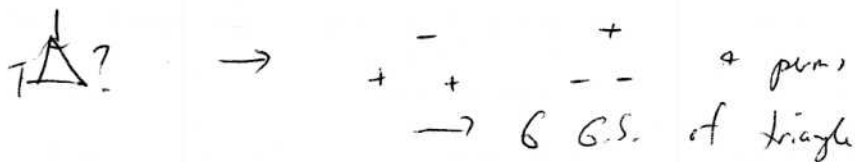
or $h_{\perp} S_i^x$ trans.

Let's ignore for now.

On bipartite lattice, H has unique G.S. up to symmetry

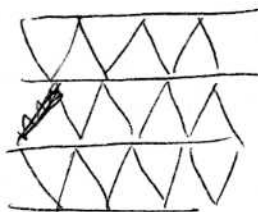


If odd-site loops are present, get frustrated.



Triangular (hexagonal) lattice

dissatisfied
one ~~frustrated~~ bond
per triangle.



removes all the frustrated bonds.

